

SUSY Working Group Summary

- SUSY \Rightarrow many free parameters:
 ↓
 large uncertainty in spectrum

different options based on \neq SUSY schemes

- SUGRA-type of models

i) minimal SUGRA:

$$m_0, M_{1/2}, A, \tan\beta, \mu \begin{pmatrix} \text{EWSB} \\ \downarrow \\ \text{sign of } \mu \\ \mu^2 \gg M_{1/2}^2 \end{pmatrix}$$

$$M_{\tilde{g}} \simeq 2.6 M_{1/2} \quad M_{\tilde{W}} \simeq 2 M_{\tilde{B}} \simeq 0.8 M_{1/2}$$

$$m_{\tilde{q}_{1,2}}^2 \simeq m_0^2 + 6 M_{1/2}^2 \quad m_{\tilde{e}_{L(R)}}^2 \simeq m_0^2 + 0.5 (0.15) M_{1/2}^2$$

$$\text{if } m_0 \gg M_{1/2} \longrightarrow m_{\tilde{q}_{1,2}} = m_{\tilde{t}} > M_{\tilde{g}}$$

$$\text{if } m_0 \ll M_{1/2} \longrightarrow M_{\tilde{g}} > m_{\tilde{q}_{1,2}} > m_{\tilde{t}} > m_{\tilde{e}} \simeq m_{\tilde{W}}$$

ii) non-universality in scalar sector

$$- m_{H_1}(M_{\text{GUT}}) \neq m_{H_2}(M_{\text{GUT}}) \neq m_0$$

$$\begin{array}{ccc} & & \\ & \swarrow & \searrow \\ m_5 & & m_{10} \end{array}$$

* may allow smaller $\mu (\ll M_{1/2})$

- Gauge Mediated SUSY models

hierarchy among masses: related to α_i ; coupl.

$$m_{\tilde{Q}}/m_{\tilde{\ell}_L} \propto \alpha_3/\alpha_2$$

$$M_{\tilde{g}} \simeq m_{\tilde{Q}} ; M_{\tilde{W}} \simeq m_{\tilde{\ell}_L} ; M_{\tilde{B}} \simeq m_{\tilde{\ell}_R}$$

Similar to SUGRA: $M_{\tilde{g}}/M_{\tilde{W}} \propto \alpha_3/\alpha_2$

$$M_{\tilde{W}}/M_{\tilde{B}} \propto \alpha_2/\alpha_1$$

large values of $\mu \Rightarrow$ gaugino-like $\tilde{\chi}_i^{\pm}$ & $\tilde{\chi}_i^0$.

- R-parity violating scenarios

$$R = (-1)^{3B + L + 2S} \left\{ \begin{array}{l} \text{SM particles: } R = 1 \\ \text{SUSY partners: } R = -1 \end{array} \right.$$

if $R \neq 1 \Rightarrow$ SUSY particles may be produced in odd number.

$$\beta_{ijk} L_i L_j \bar{E}_k + \beta'_{ijk} L_i Q_j \bar{D}_k + \cancel{\beta''_{ijk} U_i \bar{D}_j \bar{D}_k}$$

$\xrightarrow{\text{lepton # violation}}$
< bounded by present data >

easier way
to avoid proton
decay

- Flavor Mixing

S. Raby / H.-C. Cheng

generic feature of SUSY models

- if univ. scalar masses & trilinear couplings
(diag. mass matrices) at high energies
 $\Rightarrow L_e, L_\mu, L_\tau$ conserved

MSSM forbids $\mu \rightarrow e\gamma, \mu \rightarrow 3e, \mu \rightarrow e$
 $e \rightarrow \mu\pi$ conv.

- if messenger scale for SUSY is above the scale generating fermion mass structure
 \Rightarrow in general New Flavor mixing matrices appear at the gaugino-fermion-sfermion vertices

LEPTON FLAVOR VIOLATION is highly
model dependent



is a window to GUT scale Physics

One KEY Point for SUSY Scenarios

\Rightarrow light Higgs

$m_h \lesssim 130$ GeV (MSSM) fig.

$m_h \lesssim 140$ GeV (in almost all SUSY th.)

hence,

i) we expect to have seen the Higgs
at LEP: if $m_h \lesssim 100$ GeV

Tevatron: if $m_h \lesssim 120$ GeV (?)

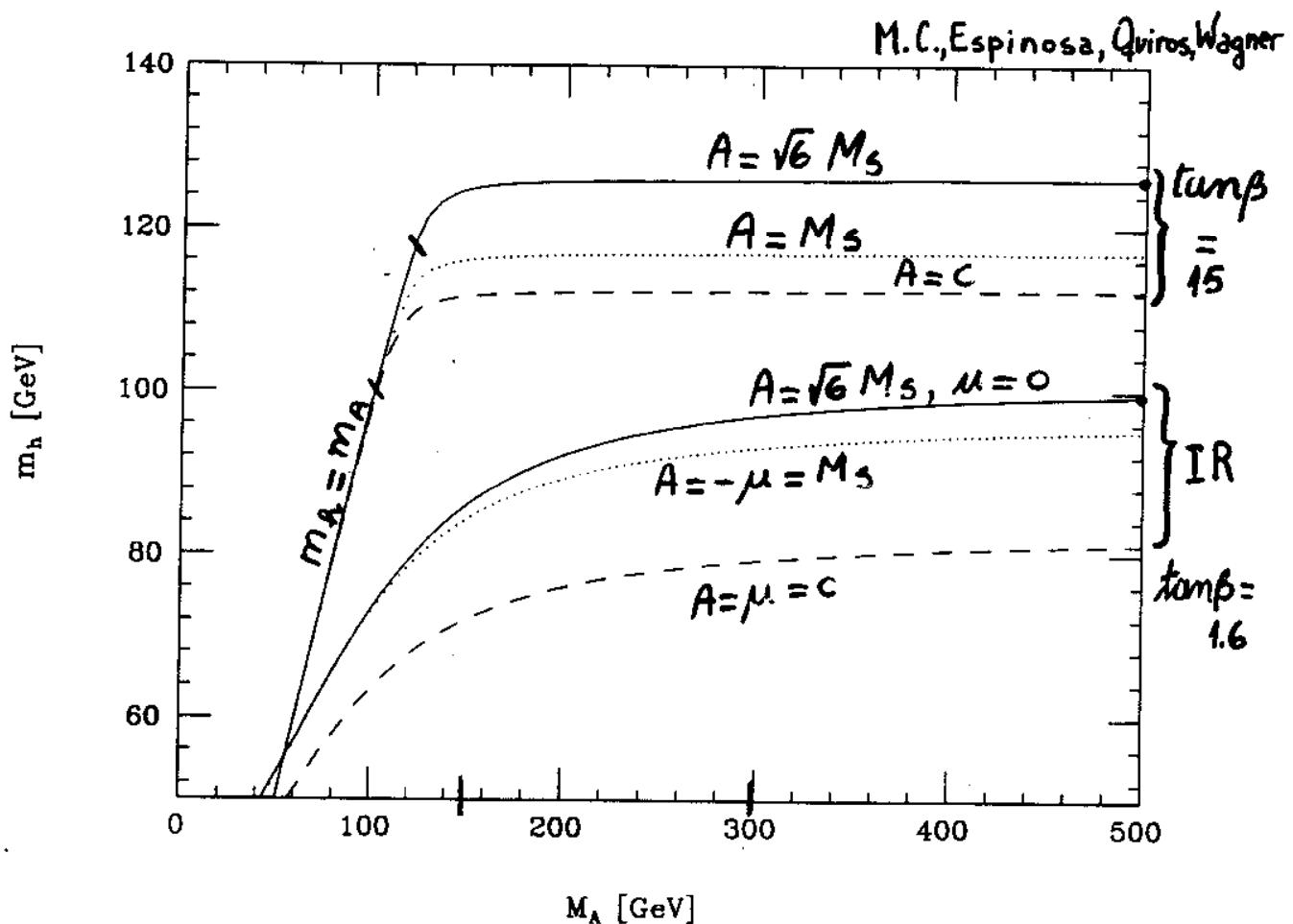
(in $W h$ or $Z h$ channels
with $25 fb^{-1}$
improved reach combining
channels — $10 fb^{-1}?$)

LHC: MSSM — discover sooner ($\lesssim 30 fb^{-1}$)
or later depending on discovery mode
NMSSM — "holes" dep. on exact spectra

ii) s-channel Higgs production at $\sqrt{s} \approx m_h$
→ first goal of the FMC see plenary talks
(to vs time issue) J. Gunion & V. Barger
+ T. Han HIG summary

Lightest Higgs Mass Vs. CP-odd Higgs Mass

- Variation with A and μ param.



- large $\tan\beta$: mixing effect very weakly dep. on μ (unless $\tan\beta \approx m_{\pm}/m_{\mu}$)
- low $\tan\beta$ (IR): important A & μ dep.

above: $M_t = 175$ GeV

$M_S = 1$ TeV

- Charginos:

Berger

$\mu^+ \mu^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-$ threshold measurement
2-point

- present LEP limit $m_{\tilde{\chi}_1^\pm} > 88 \text{ GeV}$ (close to 100 GeV after LEP)
- Tevatron reach $m_{\tilde{\chi}_1^\pm}^{\max} \simeq 200 \text{ GeV}$ if $\mathcal{L} \simeq 10 \text{ fb}^{-1}$
< good coverage for lightest chargino
but no lower limit possible >
- LHC → mainly mass difference $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$

from LEP alone \Rightarrow measurement will be for $\sqrt{s} \gtrsim 200 \text{ GeV}$

for $\sqrt{s} = 200 \text{ GeV}$ $m_{\tilde{\chi}_1^+} = 100 \text{ GeV}$

1 year $\Rightarrow 1-2 \text{ fb}^{-1} \Rightarrow \Delta m_{\tilde{\chi}_1^+} \simeq 300 \text{ MeV}$

for $\sqrt{s} = 400 \text{ GeV}$ $m_{\tilde{\chi}_1^+} = 200 \text{ GeV}$

$\mathcal{L} = 7 \text{ fb}^{-1}/\text{year}$ $\Delta m_{\tilde{\chi}_1^+} = 300 (700) \text{ MeV}$

if $m_{\tilde{\nu}_\mu} = 500 (300) \text{ GeV}$

- "comparing" with similar studies at NLC
 - ↳ Vs Energy resolution \simeq compensates
- chargino mass can be also measured by finding the endpoint in the spectrum of the chargino decay products
 - end point method \rightarrow less precision but useful to determine $\simeq m_{\tilde{\chi}^\pm}$

Polarization:

- useful tool to measure the gaugino and Higgsino components of the chargino
- only gaugino comp. couples to t-channel sneutrino exchange
 - if $\tilde{\chi}_1^+$ mainly gaugino \equiv mainly wino
 - and couples to μ_L only
 - \Rightarrow Right-handed μ can be use to turn off $\tilde{\nu}_\mu$ contribution
- turn off $\tilde{\nu}_\mu$ contribution
 <also to reduce WW background ^{better} (other ways...)>

Sleptons

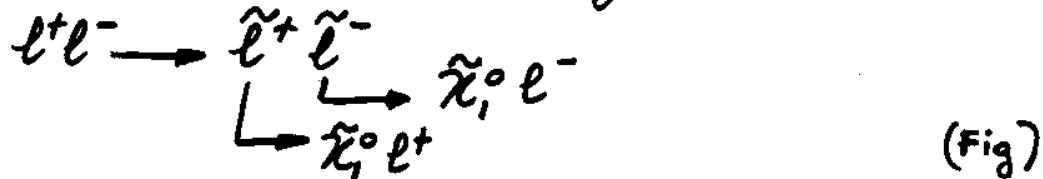
at LEP \rightarrow high \mathcal{L}_0 needed to approach kin. lim.
for $\sqrt{s} \approx 192 \text{ GeV}$ & $\mathcal{L}_0 = 500 \text{ pb}^{-1}$

$$m_{\tilde{e}} \lesssim \frac{\sqrt{s}}{2} - 5 \text{ GeV}, \quad m_{\tilde{\mu}} \lesssim \frac{\sqrt{s}}{2} - 20 \text{ GeV} \quad m_{\tilde{Z}} (?)$$

at LHC \rightarrow some masses if $\lesssim 300 \text{ GeV}$
(hard to study unless $\tilde{\chi}_1^0 \rightarrow \tilde{\ell} \ell$)

At muon collider: end point method Paige

- Select 2 acollinear leptons from



- use lepton energy distribution, find END POINTS and use that to extract: $m_{\tilde{\ell}}$ & $m_{\tilde{\chi}_1^0}$

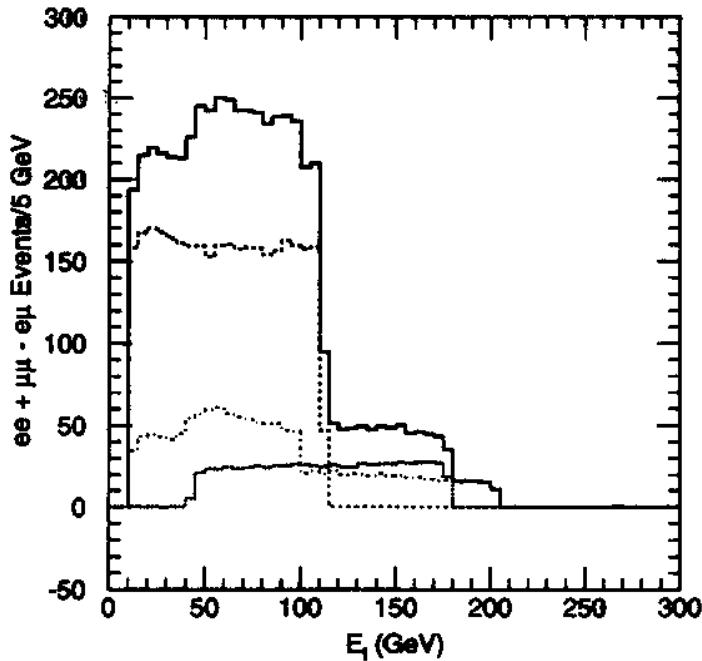
- Polarization is useful but not crucial to reject background

- Signal may be complex & require high \mathcal{L}_0
(Any FLC can complement well LHC)

Example: $m_{\tilde{e}_R} \approx 150 \text{ GeV} \Rightarrow \Delta m_{\tilde{e}_R} \approx 1 \text{ GeV}$

- Ability to measure low- E end points
is crucial! $\mathcal{L}_0 = 10 \text{ fb}^{-1}$

Signal composition: RR = dash, LL = dot, and
 RL = dash-dot:

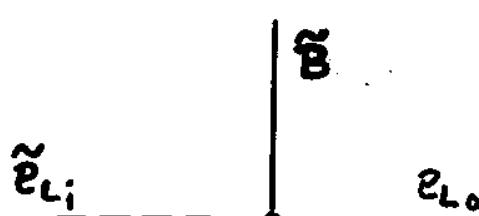


Easy to observe signal. Want to make precision measurements to complement LHC ones.

Easiest edge to measure is RR one at 111 GeV.
 Use it to measure to ~ 1 GeV

$$\frac{1}{4} \left[1 - \frac{M^2(\tilde{\chi}_1^0)}{M^2(\tilde{\ell}_R)} \right] \left[\sqrt{s} + \sqrt{s - 4M^2(\tilde{\ell}_R)} \right]$$

- New Flavor Mixings



$$W_{i\alpha} = (U_{E_L}^+ V_{E_L})_{i\alpha}$$

$$(\tilde{e}_L \tilde{\mu}_L \tilde{\chi}_L)^\ast m_{E_L}^2 \begin{pmatrix} \tilde{e}_L \\ \tilde{\mu}_L \\ \tilde{\chi}_L \end{pmatrix}$$

$$(\bar{e}_R \bar{\mu}_R \bar{\tau}_R) M_E \begin{pmatrix} e_L \\ \mu_L \\ \tau_L \end{pmatrix}$$

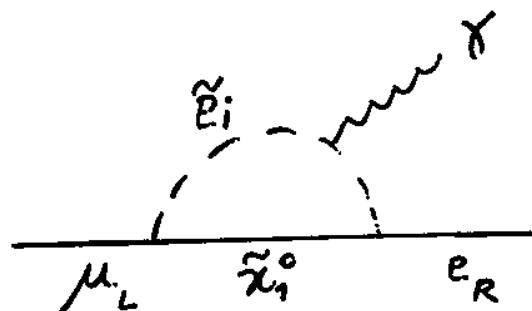
$$e_L \rightarrow V_{E_L} e_L \quad \tilde{e}_L \rightarrow U_{E_L} \tilde{e}_L$$

$$\text{in general } V_{E_L} \neq U_{E_L}$$

At muon collider

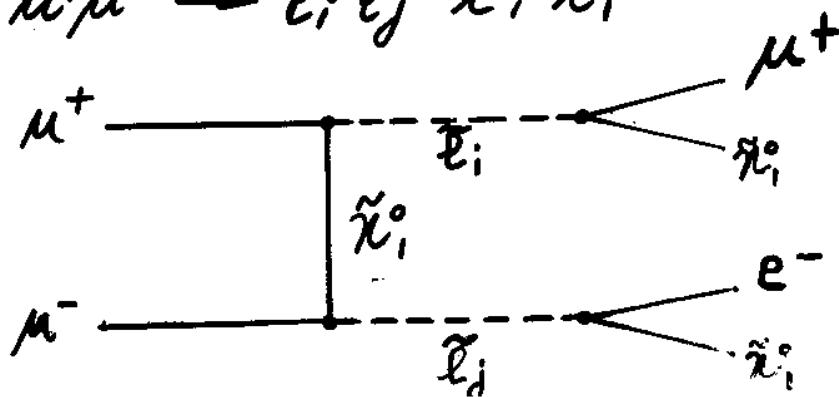
- at the Front End phase

probe by $\mu \rightarrow e \gamma$



- Direct production : flavor-changing signal

$$\mu^+ \mu^- \rightarrow \ell_i^+ \ell_j^- \tilde{\chi}_i^0 \tilde{\chi}_j^0$$

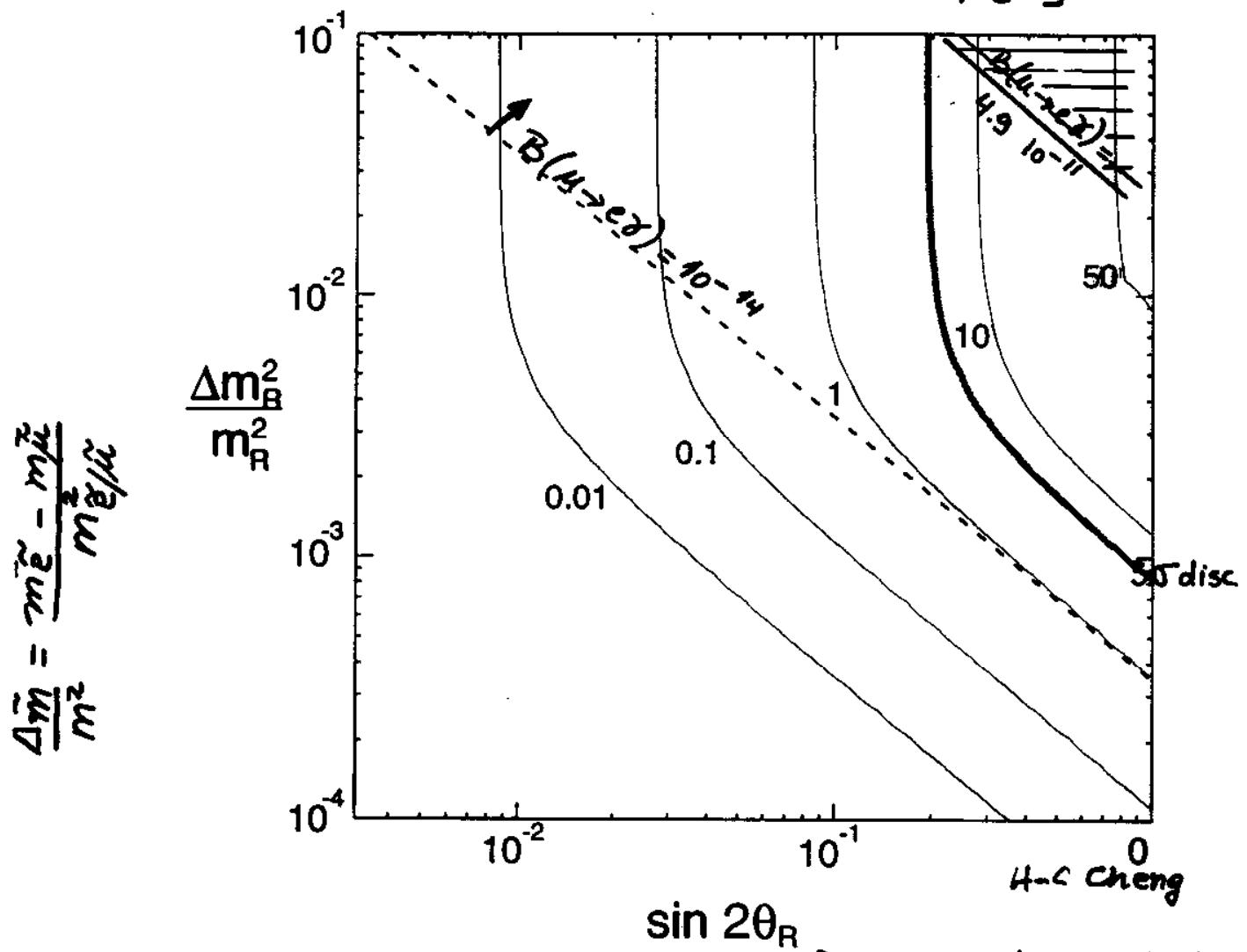


One example:

$$M_1 = 130 \text{ GeV} \quad M_2 = 260 \text{ GeV} \quad \mu = -400 \text{ GeV}$$

$$m_{\tilde{\chi}_1^0} = 200 \text{ GeV} \quad m_{\tilde{\chi}_2^0} = 350 \text{ GeV} \quad \tan\beta = 2$$

$$\Gamma(\mu^+\mu^- \rightarrow \mu^\pm e^\mp \tilde{\chi}_1^0 \tilde{\chi}_2^0) [\text{fb}]$$



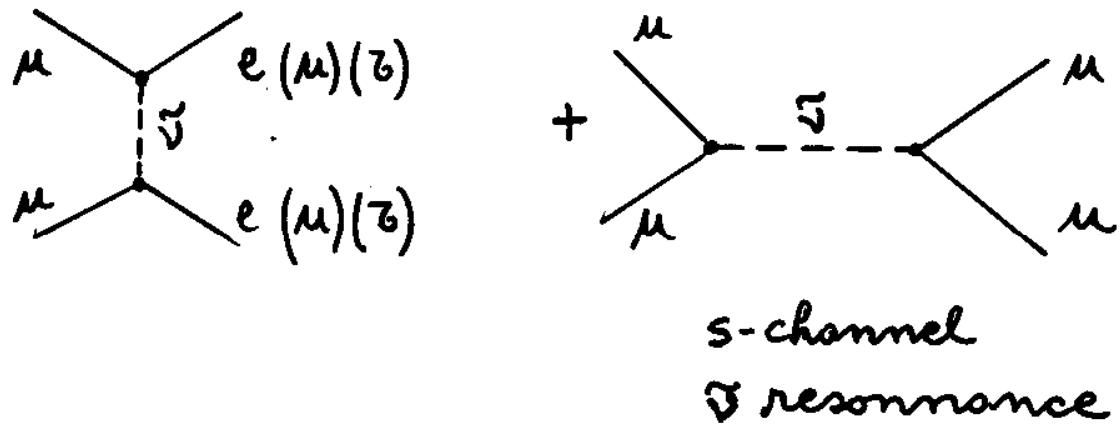
- Muon Collider & NLC complementary
($2i$ mix Vs $1i$ mix)
↓ stronger (?)

$$(\tilde{\mu} \tilde{\chi}) \begin{pmatrix} \cos\theta \sin\phi \\ -\sin\theta \cos\phi \end{pmatrix} \begin{pmatrix} \mu \\ e \end{pmatrix}_B^{\alpha}$$

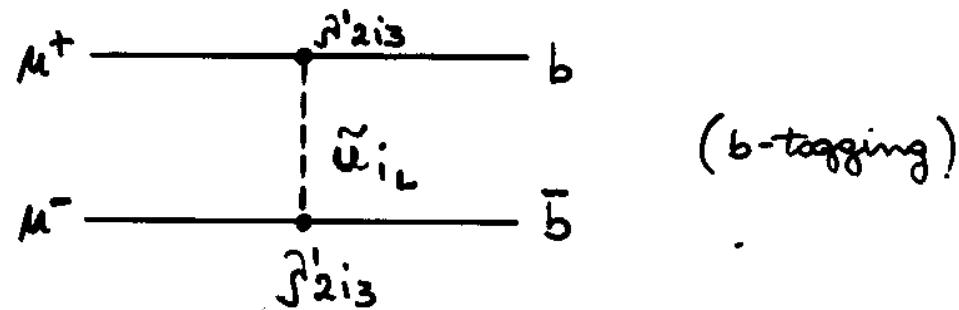
R-parity violation @ muon collider

- $\tilde{\mu}_{ijR} L_i L_j \bar{E}_R$

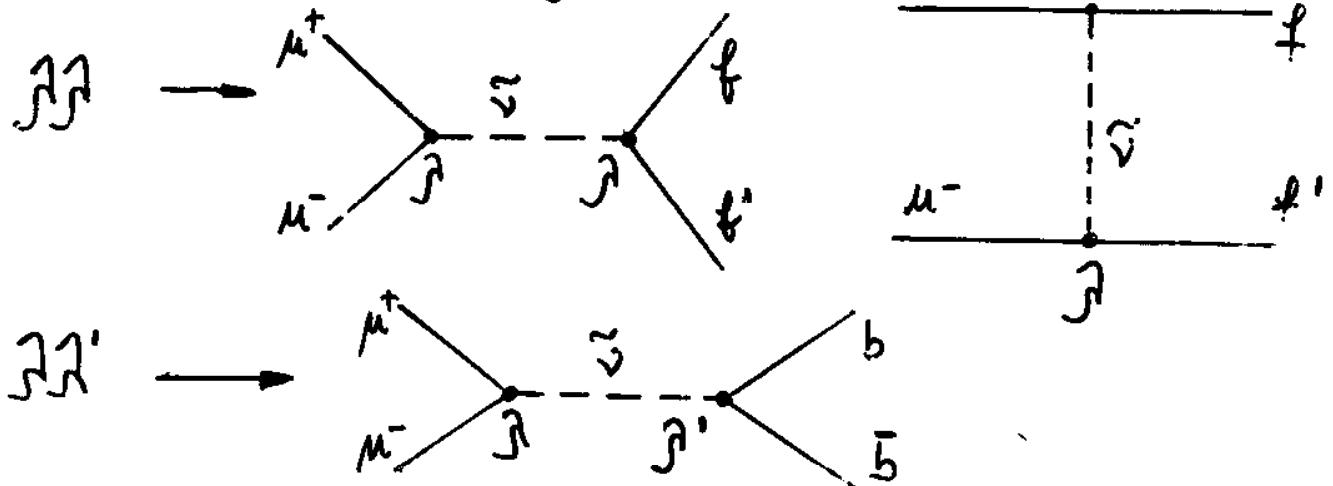
Raychaudhuri
Choudhury



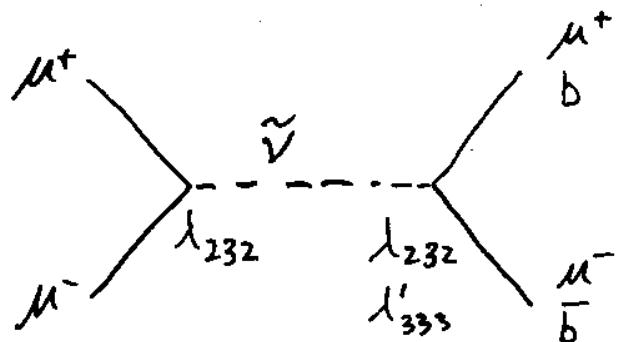
- $\tilde{\mu}'_{2i3} L_2 Q_i D_3$



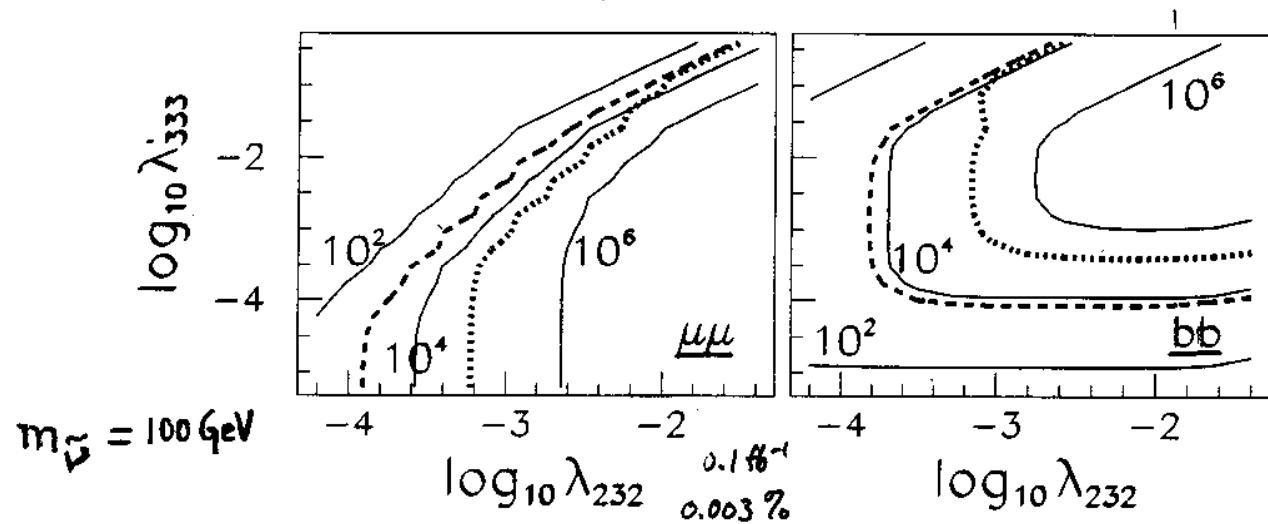
- Product of couplings



R-parity Violation $\rightarrow \tilde{\nu}$ Resonance



J. Feng talk
J. Gunion
T. Han



σ in fb : can be $O(\text{nb})$

Discover resonance for $\lambda, \lambda' \gtrsim 10^{-3}, 10^{-4}$

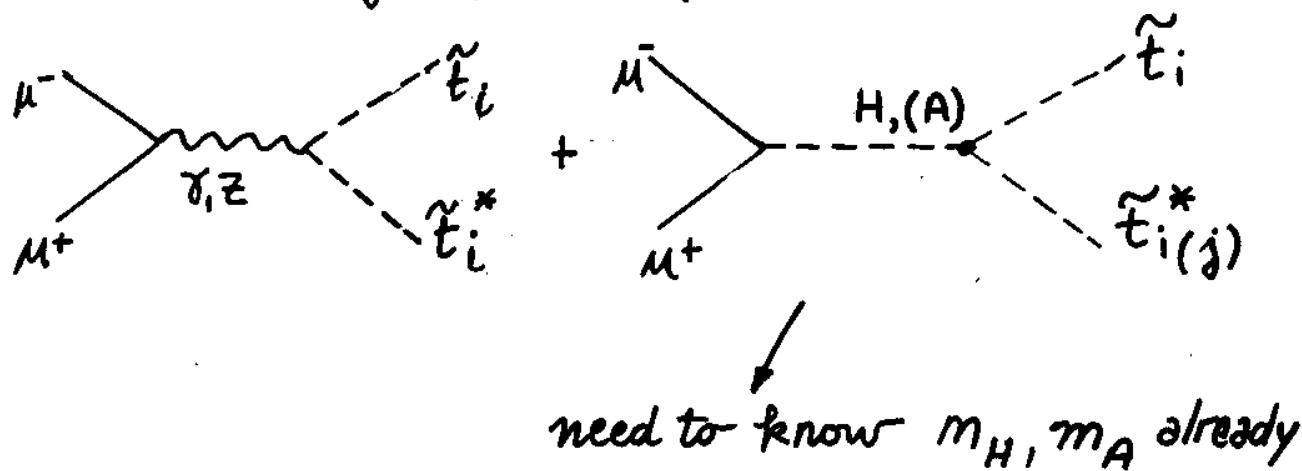
Measure R-parity couplings to $\sim 10\%$.

- \mathcal{R} t-channel only Raychaudhuri
 for $\sqrt{s} = 500 \text{ GeV}$ & 10 fb^{-1}
 can cover $\tilde{\beta}$ or $\tilde{\beta}' > 0.02$
- \mathcal{R} s-channel resonance
 $\tilde{\beta} \lesssim 10^{-2} - 10^{-3}$
 some effects within 1 year
 $\Delta \tilde{\beta} \approx 10\%$
- Complementarity of couplings which can be measured at NLC and muon Collider
 $\langle \text{LHC cannot measure any absolute value of } \tilde{\beta} \rangle$

Stop production @ muon collider

$$\mu^+ \mu^- \rightarrow \tilde{t}_i; \tilde{t}_j^* \text{ (also } \tilde{b}_i; \tilde{b}_j^*)$$

W. Porod



can be important tool to probe
neutral Higgs - stop - stop coupling
and it is sensitive to A_t/μ at
the Higgs resonance

For case considered \rightarrow fig.

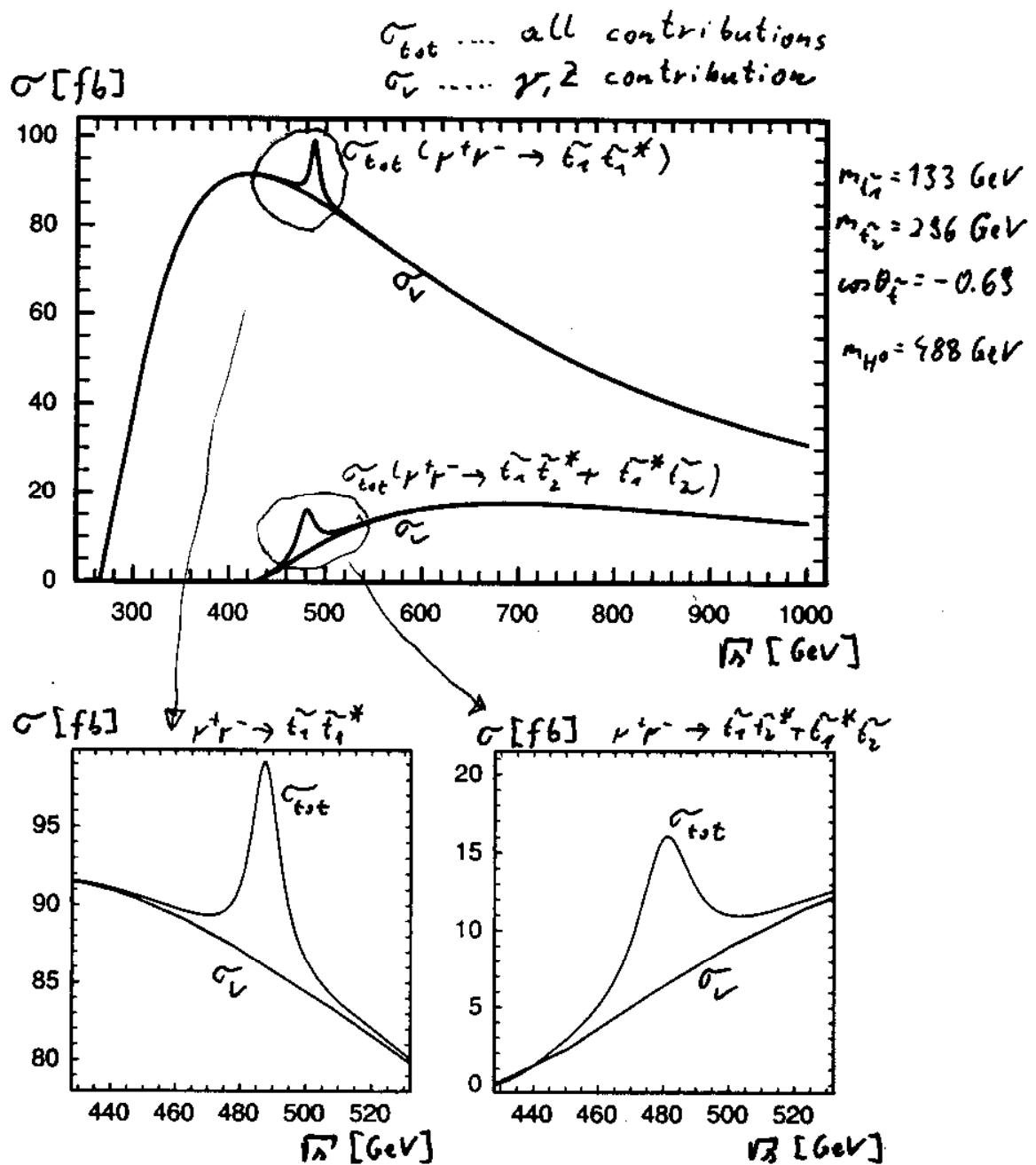
$\mu^+ \mu^- \rightarrow \tilde{t}_1 \tilde{t}_1^* \rightarrow$ 20% increase of T at resonance

$\mu^+ \mu^- \rightarrow \tilde{t}_1 \tilde{t}_2^* \rightarrow > 100\%$

preliminary studies seem to show that
at least $5-10 \text{ fb}^{-1}$ necessary $\Rightarrow H, A \gtrsim 400 \text{ GeV}$

preliminary

\tilde{t} Production



$$\begin{aligned}
 m_a &= 160 \text{ GeV}, m_u = 145 \text{ GeV}, m_d = 175 \text{ GeV}, A_t = A_b = 350 \text{ GeV} \\
 \mu &= 300 \text{ GeV}, M_2 = 740 \text{ GeV}, \tan\beta = 2, m_{A_1} = 480 \text{ GeV} \\
 m_{\tilde{t}} &= 175 \text{ GeV}, m_b = 5 \text{ GeV}
 \end{aligned}$$

Many Channel Search in Scenarios with multi-cascade decays:

SUSY Signatures/Model Discrimination at Muon Colliders

J. Kelly, UW Madison

First Muon Collider Workshop, 1997, FNAL

low energy ($\sqrt{s} = 500 \text{ GeV}$) but specially interesting for large energy ($\sqrt{s} \simeq 3.6 \text{ TeV}$)

- We calculate event numbers for pair production of all sleptons, squarks, charginos, neutralinos and Higgs bosons in the MSSM using ISAJET 7.29. Models are mSUGRA type (using ISASUGRA).
- For each pair we decay in all possible ways obtaining event numbers associated with signatures in final states described by numbers of: t, b, other jets, e^+ , e^- , μ^+ , μ^- , τ^+ , τ^- , h^0 , Z^0 , W^+ , W^- (reconstructible modes only).
- Unique signatures (with respect to other SUSY modes) can be considered a discovery tool within a given mSUGRA model; however...
- ... No SM backgrounds included in code yet. We should be wary of simple signals such as $\ell^+ \ell^- + E_{\text{miss}}$.
- In considering the mass scales in a particular scenario, we are now working on including what may be known from lower energy experiments.
- In simple two-body decays such as $\bar{e}_R^- \rightarrow e^- + \tilde{\chi}_1^0$ the energy endpoints of the electron will provide mass bounds on \bar{e}_R .
- We outline a procedure for placing bounds on 23 MSSM parameters based on the sensitivity of the signature rates (summed over produced pairs) to parameter space location. Very preliminary studies look extremely promising... *multi channel fit to constrain param. of MSSM*
- ... However, one must make sure the Monte Carlo fluctuations aren't responsible for apparent sensitivity to any of the 23 parameters.
- All of this can be accomplished without looking at final state kinematics ... just counting beans!

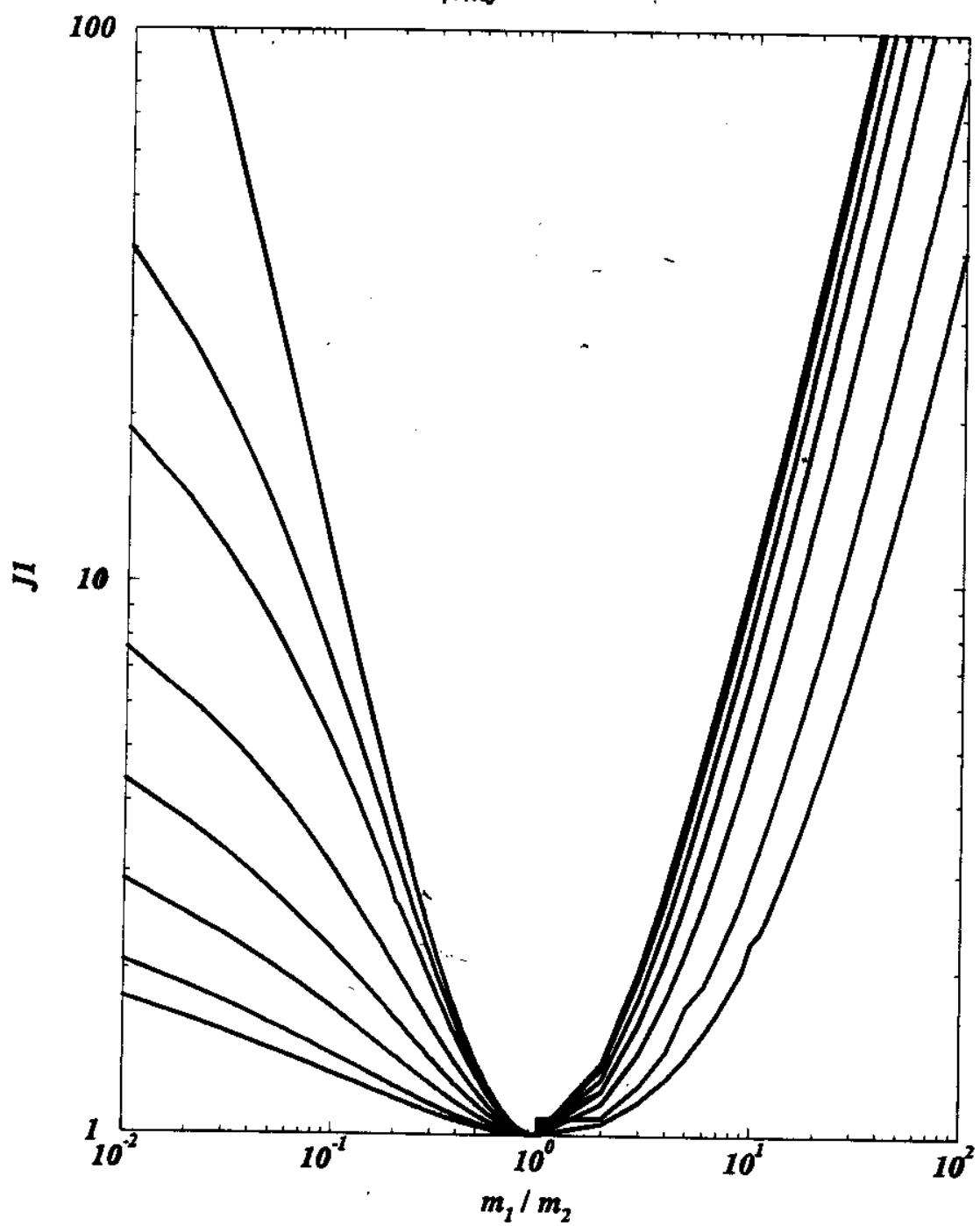
Higgs resonance



beyond minimal framework : operators at scale M_{SUSY}
Possibility of generating
radiative fermion couplings
(no Yukawa couplings at tree level)

⇒ very different (enhance) couplings
of Higgs to fermions ...

$$J_1 = \frac{\text{ACTUAL YOKO COUPLING} \rightarrow \text{EFF COUPLING}}{\text{NAIVE} \rightarrow M_B / \langle H \rangle} \quad \left(\frac{M_H}{M_B} \rightarrow 0 \right)$$



$$CP \text{ violation} \Rightarrow N(e^-\mu^+) \neq N(\mu^-e^+)$$

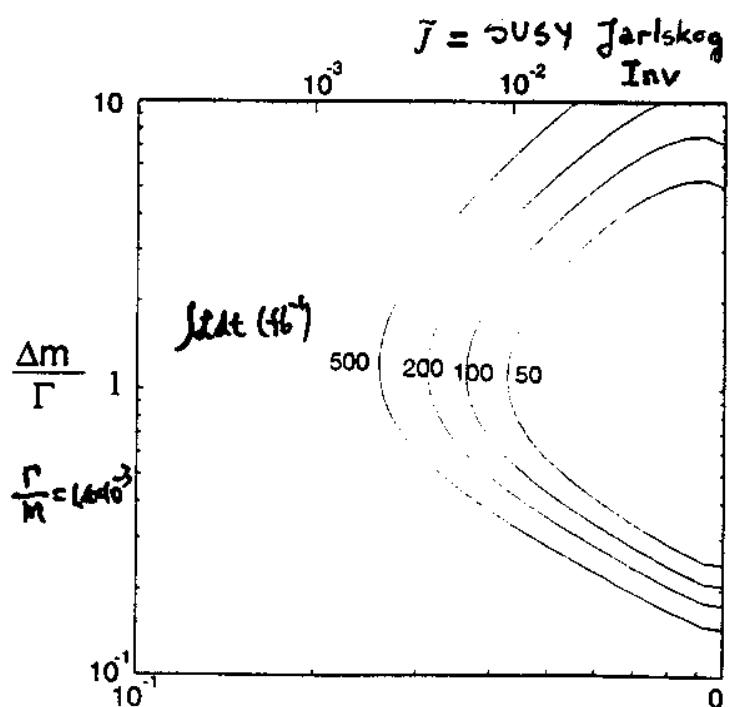
Measuring Wia provide important information about physics at high energies, e.g. GUT, flavor theories

.....

W may contain
CP phases

CP viol. signal:

asymm. between
 $N(e^+\mu^-)$ & $N(e^-\mu^+)$



$$\frac{\Delta m}{\Gamma} = 10^{-3}$$

$$M_1 = 100 \text{ GeV}$$

$$M_2 = 200 \text{ GeV}$$

$$\mu = -400 \text{ GeV}$$

$$tan\beta = 2$$

$$M_{\tilde{\chi}_1^0} = 150 \text{ GeV}$$

$$\sin 2\theta$$

$$\theta_u = \theta_d = \theta_s = \theta$$

$$\sin \delta = 1$$

After LEP, Tevatron & LHC

for SUSY measurements we need

the FUTURE Lepton Colliders

NLC, JLC, DESY LC and/or MC

→ better chances for

- heavy Higgs $A, H, H^\pm \rightarrow$ specially if SUSY decays open.
- sleptons
- in general for precision measurements of masses, couplings, spin, widths
chargino mixing, flavor mixing

NLC-type: large electron long. polarization (+ high \mathcal{L}_0)
useful to enhance certain signals &
kill backgrounds + probe $H\tilde{W}$ mixing in $\tilde{\chi}$;

MC → highest energy reach ($\approx 3\text{TeV}$) with
gradual upgrade involving MANY competitive
precision measurements studies

- superior beam energy resol. (trade with $\mathcal{L}_0 \dots$)
- s channel Higgs factory, $\tilde{\nu}$ factory ($\tilde{\chi}$)

+ NLC: complement. in probing new flavor mix. & $\tilde{\chi}$ couplings